

## Dynamical Systems II

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### Coursework 1 (Solutions)

Hand in the complete solutions to all three questions in the general office (room C123).

DEADLINE: Tuesday 10/11/2009 at 16:00

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1) Consider the linear dynamical system of the form

$$\begin{aligned} \dot{x}_1 &= ax_1 + bx_2 \\ \dot{x}_2 &= cx_1 + dx_2 \end{aligned} \quad \text{with } a, b, c, d \in \mathbb{R}. \quad (1)$$

The Jacobian is

$$J = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

The eigenvalues are computed from  $\det(J - \lambda\mathbb{I}) = 0$ , i.e.

$$\lambda^2 - (a + d)\lambda + ad - bc = 0,$$

which is solved by

$$\lambda_{\pm} = \frac{a + d}{2} \pm \frac{1}{2} \sqrt{a^2 + 4bc - 2ad + d^2} = \alpha \pm i\beta.$$

This means

$$a^2 + 4bc - 2ad + d^2 = (a - d)^2 + 4bc < 0. \quad (2)$$

Let us now assume that the equations (1) have a solution of the general form

$$x_1 = Ae^{\lambda t}, \quad x_2 = Be^{\lambda t} \quad \text{for } A, B \in \mathbb{R}. \quad (3)$$

Then, substituting this into (1) gives

$$\begin{aligned} A\lambda e^{\lambda t} &= aAe^{\lambda t} + bBe^{\lambda t}, \\ B\lambda e^{\lambda t} &= cAe^{\lambda t} + dB e^{\lambda t}, \end{aligned}$$

which corresponds to

$$\begin{aligned} (a - \lambda)A + bB &= 0, \\ cA + (d - \lambda)B &= 0, \end{aligned} \quad \Leftrightarrow \quad (J - \lambda\mathbb{I}) \begin{pmatrix} A \\ B \end{pmatrix} = 0.$$

These equations only have nontrivial solutions when  $\det(J - \lambda \mathbb{I}) = 0$ , such that the general solution (3) becomes

$$\begin{aligned} x_1 &= A_1 e^{\lambda+t} + A_2 e^{\lambda-t} = e^{\alpha t} (A_1 e^{i\beta t} + A_2 e^{i\beta t}), \\ x_2 &= B_1 e^{\lambda+t} + B_2 e^{\lambda-t} = e^{\alpha t} (B_1 e^{-i\beta t} + B_2 e^{-i\beta t}). \end{aligned}$$

Clearly for  $\alpha < 0$  we have  $x_{1/2} \rightarrow 0$  for  $t \rightarrow \infty$ .

Next we establish that the trajectory spirals into the origin. Using  $x_1 = r \cos \vartheta$  and  $x_2 = r \sin \vartheta$  we have  $\vartheta = \arctan x_2/x_1$ , such that

$$\begin{aligned} \dot{\vartheta} &= \frac{\partial \vartheta}{\partial x_1} \dot{x}_1 + \frac{\partial \vartheta}{\partial x_2} \dot{x}_2 = -\frac{x_2 \dot{x}_1}{x_1^2 + x_2^2} + \frac{x_1 \dot{x}_2}{x_1^2 + x_2^2} \\ &= \frac{x_1(cx_1 + dx_2) - x_2(ax_1 + bx_2)}{x_1^2 + x_2^2} = \frac{cx_1^2 - bx_2^2 + (d-a)x_1x_2}{x_1^2 + x_2^2}. \end{aligned}$$

We have used here  $\partial \vartheta / \partial x_1 = -x_2 / (x_1^2 + x_2^2)$  and  $\partial \vartheta / \partial x_2 = x_1 / (x_1^2 + x_2^2)$ .

From (2) we know that  $b$  and  $c$  must have opposite signs.

- for  $x_2 = 0$  we have  $\dot{\vartheta} = c > 0$  for  $b < 0, c > 0$  or  $\dot{\vartheta} = c < 0$  for  $b > 0, c < 0$
- for  $x_2 \neq 0$  we have  $\dot{\vartheta} \neq 0$  or  $cx_1^2 - bx_2^2 + (d-a)x_1x_2 = 0$ . From the latter follows

$$c \frac{x_1^2}{x_2^2} - b + (d-a) \frac{x_1}{x_2} = 0,$$

which means

$$\frac{x_1}{x_2} = \frac{a-d}{2c} \pm \frac{1}{2c} \sqrt{a^2 + 4bc - 2ad + d^2}. \quad (4)$$

As we know from (3) that  $a^2 + 4bc - 2ad + d^2 < 0$  there is no real solution to (4) and therefore  $\dot{\vartheta} \neq 0$  for  $x_2 \neq 0$ .

Therefore when  $b < 0, c > 0, a + d < 0$  the trajectories spiral anti-clockwise into the origin and when  $b > 0, c < 0, a + d < 0$  the trajectories spiral clockwise into the origin.

$\Sigma = 20$

2) We have

$$\begin{aligned} \dot{x}_1 &= x_2 + x_1(3 - x_1^2 - x_2^2)(x_1^4 + x_2^4 + 2x_1^2x_2^2 - 1)^2 \\ \dot{x}_2 &= -x_1 + x_2(3 - x_1^2 - x_2^2)(x_1^4 + x_2^4 + 2x_1^2x_2^2 - 1)^2 \end{aligned}$$

i) With  $x_1 = r \cos \vartheta$  and  $x_2 = r \sin \vartheta$ , we obtain

$$\dot{x}_1 = \dot{r} \cos \vartheta - r \sin \vartheta \dot{\vartheta} = r \sin \vartheta + r \cos \vartheta (3 - r^2)(r^4 - 1)^2 \quad (5)$$

$$\dot{x}_2 = \dot{r} \sin \vartheta + r \cos \vartheta \dot{\vartheta} = -r \cos \vartheta + r \sin \vartheta (3 - r^2)(r^4 - 1)^2 \quad (6)$$

where we used  $x_1^2 + x_2^2 = r^2$  and  $x_1^4 + x_2^4 + 2x_1^2x_2^2 = r^4$ .

Computing (5)cos  $\vartheta$  + (6)sin  $\vartheta$  gives

$$\dot{r} = r(3 - r^2)(r^4 - 1)^2 \quad (7)$$

and  $(6)\cos\vartheta-(5)\sin\vartheta$  yields

$$r\dot{\vartheta} = -r \Rightarrow \dot{\vartheta} = -1. \tag{8}$$

Since  $\dot{\vartheta} \neq 0$  for all time, the origin is the only fixed point. 5

ii) With (7) we compute  $\dot{r}$  at the inner and outer boundary of

$$\mathcal{D} = \{(r, \vartheta) : 1/2 \leq r \leq 2\} .$$

We find

$$\begin{aligned} \dot{r}(r = 1/2) &= \frac{1}{2} \frac{11}{4} \left(-\frac{15}{16}\right)^2 > 0 \\ \dot{r}(r = 2) &= 2(-1)15^2 < 0. \end{aligned}$$

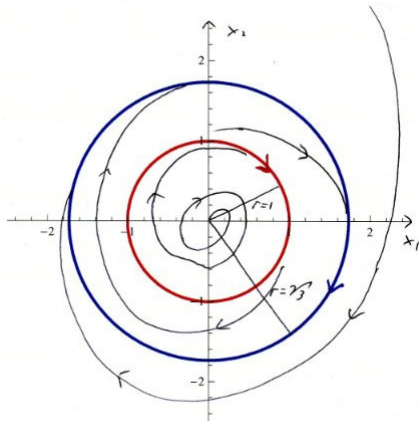
$\Rightarrow$  trajectories which enter  $\mathcal{D}$  never leave it.

$\Rightarrow$  It follows by the Poincaré-Bendixson theorem that there is at least one limit cycle in  $\mathcal{D}$ . 3

iii) We have

$$\begin{aligned} \dot{r} &= 0 && \text{for } r = 0, 1, \sqrt{3} \equiv \text{limit cycles} \\ \dot{r} &> 0 && \text{for } 0 < r < 1 \\ \dot{r} &> 0 && \text{for } 1 < r < \sqrt{3} \\ \dot{r} &< 0 && \text{for } \sqrt{3} < r \\ \dot{\vartheta} &= -1 && \equiv \text{clockwise direction} \end{aligned}$$

Assembling all this gives the following phase portrait. 5



iv) The  $\alpha$  and  $\omega$  limit sets are

$$L_\alpha(\vec{x}) = \begin{cases} (0, 0) & \text{for } r < 1 \\ C_1 & \text{for } 1 \leq r < \sqrt{3} \\ C_{\sqrt{3}} & \text{for } r = \sqrt{3} \\ \emptyset & \text{for } \sqrt{3} < r \end{cases} \quad L_\omega(\vec{x}) = \begin{cases} (0, 0) & \text{for } r = 0 \\ C_1 & \text{for } 0 < r \leq 1 \\ C_{\sqrt{3}} & \text{for } 1 < r \end{cases}$$

- $C_1$  is semistable and  $C_{\sqrt{3}}$  is stable.
- These findings are consistent with the conclusions drawn in *ii*), both limit cycles are in  $\mathcal{D}$ .
- Defining a domain

$$\mathcal{D}' = \left\{ (r, \vartheta) : \sqrt{3} - \varepsilon \leq r \leq \sqrt{3} + \varepsilon \right\}$$

we can use the same arguments as in *ii*) to conclude that there is a limit cycle in  $\mathcal{D}'$ .

- However, we can not define such an annular region around  $C_1$ , as for

$$\mathcal{D}'' = \{(r, \vartheta) : 1 - \varepsilon \leq r \leq 1 + \varepsilon\}$$

the trajectories will enter  $\mathcal{D}''$  on the inner boundary and leave it at the outer boundary. Thus the Poincaré-Bendixson theorem can not be applied in this case.

- v) Bendixson's criterium requires the region to be simply connected, which is not the case for  $\hat{\mathcal{D}}$ . Therefore the criterium can not be applied.

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$\Sigma = 20$

3) We have

$$\begin{aligned} \dot{x}_1 &= x_1 x_2^2 - 9x_1 - 16x_2^3 \\ \dot{x}_2 &= 4x_1 x_2^2 + 2x_2 x_1^2 \end{aligned}$$

i) **Lyapunov stability theorem:** Consider the system  $\dot{\vec{x}} = \vec{X}(\vec{x})$  with a fixed point at the origin. If there exists a real valued function  $V(\vec{x})$  in a neighbourhood  $N(\vec{x} = 0)$  such that:

- the partial derivatives  $\partial V/\partial x_1, \partial V/\partial x_2$  exist and are continuous
- the function  $V(\vec{x})$  is positive definite
- $dV/dt$  is negative semi-definite (definite)

then the origin is a stable (asymptotically stable) fixed point.

**Def.:** A function  $V$  for which the conditions i)-iii) hold with iii) semi-definite is called weak Lyapunov function.

**Def.:** A function  $V$  for which the conditions i)-iii) hold with iii) definite is called strong Lyapunov function.

Verify the requirements a)-c) for  $V(x_1, x_2) = 4x_1^2 + 16x_2^2$

- Clearly the partial derivatives  $\partial V/\partial x_1, \partial V/\partial x_2$  exist and are continuous.
- $\because V(0, 0) = 0$  and  $V(x_1, x_2) > 0$  for  $(x_1, x_2) \neq 0$   
 $\Rightarrow$  the function  $V(x_1, x_2) = 4x_1^2 + 16x_2^2$  is positive definite

c) Compute  $dV/dt$ :

$$\begin{aligned}\dot{V} &= \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 \\ &= 8x_1(x_1x_2^2 - 9x_1 - 16x_2^3) + 32x_2(4x_1x_2^2 + 2x_2x_1^2) \\ &= 8x_1^2x_2^2 - 72x_1^2 - 128x_1x_2^3 + 128x_1x_2^3 + 64x_2^2x_1^2 \\ &= -8x_1^2[9 - x_2^2(1 + 8)] \\ &= -72x_1^2(1 - x_2^2)\end{aligned}$$

- $\Rightarrow \dot{V} = 0$  for  $(0, x_2)$
- $\Rightarrow \dot{V} \leq 0$  for  $|x_2| < 1$
- $\Rightarrow$  is negative semi-definite

$\Rightarrow V$  is a weak Lyapunov function

$\Rightarrow$  by the Lyapunov stability theorem follows that the origin is a stable fixed point. □6

ii) For  $|x_2| < 1$  all points inside the level curve of  $V(x_1, x_2) = 4x_1^2 + 16x_2^2$  will be dragged to the origin.

$$\Rightarrow 4x_1^2 + 16x_2^2 < 16 \quad \Rightarrow \frac{x_1^2}{4} + x_2^2 < 1$$

$\Rightarrow$  the length of the minor is 2 and the length of the major is 4. □2

iii) **Corollary:** Let  $V(\vec{x}(t))$  be a weak Lyapunov function for the system  $\dot{\vec{x}} = \vec{X}(\vec{x})$  in a neighbourhood of an isolated fixed point  $\vec{x}_f = (0, 0)$ . Then if  $\dot{V} \neq 0$  on a trajectory, except for the fixed point, the origin is asymptotically stable.

For  $(0, x_2)$  the dynamical system reduces to  $\dot{x}_1 = -16x_2^3$  and  $\dot{x}_2 = 0$ . □2

This means the line  $(0, x_2)$  is not a trajectory and therefore it follows from the corollary that the origin is asymptotically stable. □ $\Sigma = 10$