

Complex Spectrum of a Spontaneously Unbroken PT Symmetric Hamiltonian



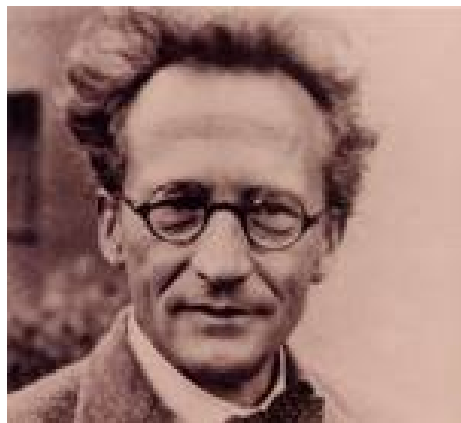
Anadolu University

Cem Yuce
17th of July 2007

Outline

1. A spontaneously broken Hamiltonian with real spectra
2. Symmetry Property
3. A new approach
4. Conclusions

Hamiltonians must
be Hermitian



~1930

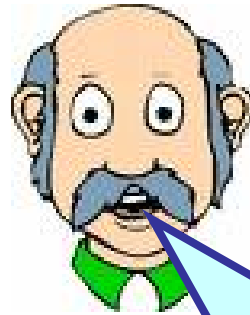
should be replaced
by the PT symmetry

PRL 80 5243
(1998)

1998



PT Symmetry condition



**there
exist also examples with
real spectra for which
Hamiltonian is not PT
symmetric.**



PT Symmetry

$$p \rightarrow p, \quad x \rightarrow -x, \quad i \rightarrow -i, \quad t \rightarrow -t$$

$$[H, PT] = 0 ; \quad PT\Psi(x) = \mp \Psi(x)$$



**Spontaneous breaking of PT transformations:
The appearances of complex eigenvalues**

A Time-dependent Hamiltonian

$$H = p^2 + x^2 + 2i f(t) x$$

PT Symmetric: $f(-t) = f(t)$

Analytic Solution

$$\Psi_n = \exp\left(-i\Phi(t) + \alpha z - \frac{z^2}{2}\right) H_n(z)$$

$$z = x + ig(t)$$

$$\dot{\alpha} = 2(f - g)$$

$$\alpha = \frac{\dot{g}}{2}$$

For $n=0, n=1$

$$H_0(x + ig) = 1$$

$$H_1(x + ig) = 2(x + ig)$$

C. Yuce, [quant-ph/0703235]

“

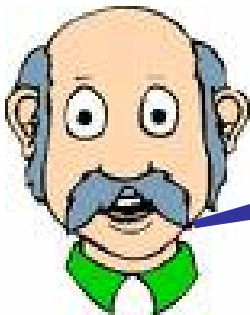
Energy Spectrum

for $n=0$

$$\langle 0|E|0 \rangle = \frac{\int_{-\infty}^{\infty} \Psi_0^* H \Psi_0 dx}{\int_{-\infty}^{\infty} |\Psi_0|^2 dx} = 1 + g^2 + \alpha^2 + 2i\alpha f$$

for $n=1$

$$\langle 1|E|1 \rangle = \frac{E_{11} + 2i\alpha f(3 + 2\alpha^2 + 2g^2)}{1 + 2g^2 + 2\alpha^2}$$



Spectrum is not real unless $f \alpha=0$

Suppose

$$f(t) = f_0 \Rightarrow \alpha(t) = 0 \quad (H = p^2 + x^2 + 2if_0x)$$

Examples

$$V(x) = x^2 + 2i t^2 x$$

PT Symmetric

PT Symmetric

$$\Psi = \exp \left(t \left(x + i \left(t^2 - \frac{1}{2} \right) \right) - \frac{1}{2} \left(x + i \left(t^2 - \frac{1}{2} \right) \right)^2 \right) \\ \times e^{-2i(n+\frac{1}{2})t - i \left(\frac{t^5}{5} + \frac{t^3}{3} - \frac{t}{4} \right)} H_n \left(x + i \left(t^2 - \frac{1}{2} \right) \right).$$

$$\langle 0 | E | 0 \rangle = \frac{5}{4} + t^4 + 2it^3$$



Outline

1. A spontaneously broken Hamiltonian with real spectra
2. Symmetry Property
3. A new approach
4. Conclusions

Symmetry Property

Theorem:

$$\text{Let } H = \frac{p^2}{2m} + U^R + iU^I$$

The spectrum is real if $\langle U^I \rangle = \int |\Psi|^2 U^I d^3x = 0$

Hermiticity condition

1 $U^I = 0$

generalization of PT condition

2 $U^I(-x, t) = -U^I(x, t);$
 $|\Psi(-x, t)|^2 = |\Psi(x, t)|^2$



**PT condition works only for
time-independent potentials**

$$U^I(-x, t) = -U^I(x, t) ;$$
$$| \Psi(-x, t) |^2 = | \Psi(x, t) |^2$$



But, they work for both cases.

$$V(x) = x^2 + 2i t^2 x$$

PT Symmetric

PT Symmetric

$$\Psi = \exp \left(t \left(x + i \left(t^2 - \frac{1}{2} \right) \right) - \frac{1}{2} \left(x + i \left(t^2 - \frac{1}{2} \right) \right)^2 \right) \\ \times e^{-2i(n+\frac{1}{2})t - i \left(\frac{t^5}{5} + \frac{t^3}{3} - \frac{t}{4} \right)} H_n \left(x + i \left(t^2 - \frac{1}{2} \right) \right).$$

$$U^I(-x, t) = -U^I(x, t) ;$$

$$| \Psi(-x, t) |^2 \neq | \Psi(x, t) |^2$$

Outline

1. A spontaneously broken Hamiltonian with real spectra
2. Symmetry Property
3. A new approach
4. Conclusions

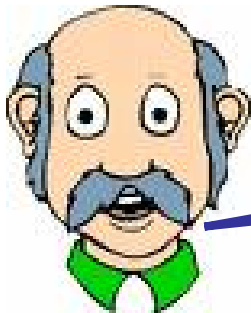
$$H = H_0 + i\lambda V$$

$$V_{nk} = \langle n|V|k \rangle$$

The energy shift:

$$\Delta_n = i \lambda V_{nn} - \lambda^2 \sum_{k \neq n} \frac{|V_{nk}|^2}{E_n^0 - E_k^0} - i \lambda^3$$

$$\left(\sum_{k \neq n} \sum_{l \neq n} \frac{V_{nk} V_{kl} V_{ln}}{(E_n^0 - E_k^0)(E_n^0 - E_l^0)} - \sum_{k \neq n} \frac{|V_{nk}|^2 V_{nn}}{(E_n^0 - E_k^0)^2} \right) + \dots$$



**Spectrum is real if
all of the complex terms vanish!!**

- $V_{nn} = 0$
- $V_{nk}V_{kl}V_{ln} = 0$
- $V_{nk}V_{kl}V_{ls}V_{sp}V_{pn} = 0$
- ...



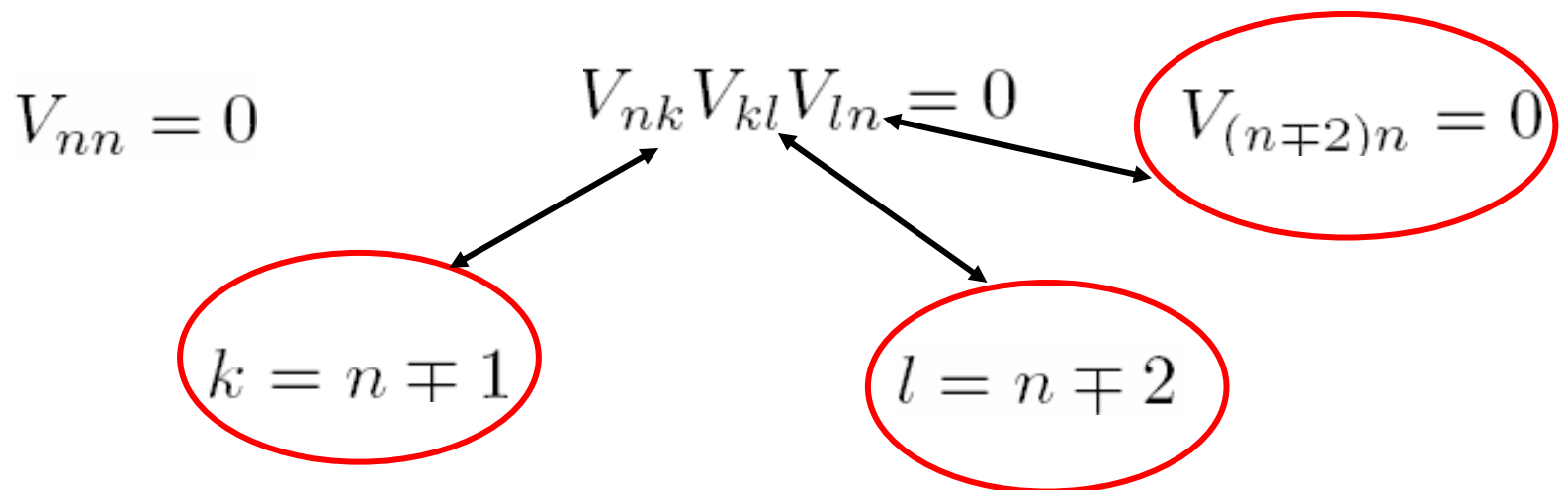
But, there are infinitely many terms.

Let's use a and a+ operators.

Examples

$$H = p^2 + x^2 + i\lambda x \quad x = \frac{1}{\sqrt{2}} (\hat{a}^\dagger + \hat{a})$$

$$\langle m|x|n \rangle = \frac{1}{\sqrt{2}} (\sqrt{n} \delta_{m,n-1} + \sqrt{n+1} \delta_{m,n+1})$$



Examples

$$H = p^2 + x^2 + i\lambda x^3$$

$$\langle m | x^3 | n \rangle \neq 0 \quad \text{if } m \mp 1 \text{ or } m = \mp 3$$

$$V_{nn} = 0$$

$$V_{nk} V_{kl} V_{ln} = 0$$

$$V_{(n \mp 2)n} = 0$$

$$V_{(n \mp 4)n} = 0$$

$$k = n \mp 1$$
$$k = n \mp 3$$

$$l = n \mp 2$$
$$l = n \mp 4$$

Examples

$$iV = i\lambda f(\hat{a}) \qquad iV(x) = i\lambda f(\hat{a}^\dagger)$$

$$iV = i\lambda e^{\hat{a}^\dagger} = i\lambda e^{(x-ip)}$$

$$iV = i\lambda \hat{a}^\dagger \hat{a} \hat{a}^\dagger$$

Outline

1. A spontaneously broken Hamiltonian with real spectra
2. Symmetry Property
3. A new approach
4. Conclusions

Summary of results

- Spontaneously unbroken PT symmetric Hamiltonians may have real spectra.
- New approach: Perturbation theory is useful for searching for a non-Hermitian Hamiltonian with real spectrum.